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A new Approach for Scaling Up from a Small Cold to a Large Hot Bed for Biomass Conversion in a Bubbling Fluidized Bed Reactor

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Abstract

Bubbling fluidized beds are simple and attractive means of achieving efficient conversion of biomass if particle segregation and the associated effects are minimized. To improve the knowledge of fluidized bed reactor design, this paper compares the behavior of a hot bed containing a certain amount of biomass with the behavior in a cold bed having the same biomass loads and particle properties. An approach for scaling up a cold bed to a large hot bed for the same volume fraction of biomass is introduced. The proposed scheme uses the bed expansion ratio as an output from the cold bed. This approach provides an accurate means of attaining dynamic similarity in bubbling behavior between two different beds without constraining the fluid and particle properties as well as the bed height.

Keywords: Biomass; Scale up; Bed expansion; CFPD; Bubbling fluidized bed

1 Introduction

To enhance the understanding of fluidized bed behavior for reactor designs and scale up, this paper compares the bubbling behavior in a hot fluidized bed containing biomass with that observed in Agu *et al* (2019a) under cold flow condition using the same bed material and biomass properties. In the literature (Agu *et al*, 2019a), the mixing and segregation pattern of wood chips and pellets of a wide density difference but similar volume-equivalent spherical particle diameter was studied using electrical capacitance tomography (ECT).

Understanding how the cold bed behavior is related to the behavior in an active hot bed is often a challenge in modelling and scaling up. The traditional scaling up approaches (Glicksman *et al*, 1993; Glicksman, 1988; Glicksman, 1984; Romero and Johanson, 1962) that have been used in many studies (Kraft *et al*, 2018; Biglari *et al*, 2016; Thapa and Halvorsen, 2014) do not consider the influence of feedstock and volatiles (i.e. biomass) on the bubbling behavior. As shown in different studies (Agu *et al*, 2019a; Fotovat *et al*, 2015a), the amount of biomass in a bed affects the bubble diameter, bubble frequency and bed expansion, and also

causes a delay in slug flows. In an attempt to keep the ratios of diameter and density of fuel particles to those of a bed material constant, the two beds will have different fluid and particles properties. Different scaling up approaches accounting for intrinsic behavior in a bed have also been demonstrated. On the account of segregation effects, a scaling approach based on the biomass residence time distribution was reported in (Zhang and Xu, 2015). Scaling up a bed based on the segregation effects on the thermochemical conversion of biomass can also be achieved by using the relative time scales (Damköhler numbers) between the segregation and devolatilization of the fuel particles (Salatino and Solimene, 2017). Moreover, a bubbling fluidized bed can also be scaled up by keeping at least one of the bubble properties constant (Agu *et al*, 2019b).

In this study, a new approach for scaling a bubbling bed behavior from a cold small to a hot large bed is introduced based on a combination of Horio *et al* (1986) scaling law and correlations for minimum fluidization velocity and bed expansion ratio. The scaling approach introduced allows the same properties of biomass as well as the same bed material type to be used in both beds. For the scaled bed, the Wen and Yu (1966a) correlation is used to determine the particle size while the model for bed expansion ratio (Agu *et al*, 2019c) is used to determine the aspect ratio (bed height to diameter) owing to the fact that bubble growth and bed expansion depend on the bed aspect ratio. The scaled bed aspect ratio is also verified by simulation of the two different beds using the Barracuda software based on the CFPD (computational particle-fluid dynamics) model.

2 Modelling and scaling up

To scale up fluidized bed behavior from a small to a larger bed diameter, scaled and transferrable information from the smaller (reference) bed is required. In a hot flow bed, temperature and pressure are easier to measure. Temperature is not a product of bubble flows but the distribution of temperature as well as the pressure fluctuations can be linked to the bubbling behavior. However, modelling of temperature distribution and pressure fluctuation is a complex task.

In a cold bed, different properties can easily be measured, including bubble diameter, bubble frequency and bed expansion. Different correlations are also available in literature for prediction of the different bubble properties for a bed of single particle type, e.g. sand particles. Based on the results reported in Agu *et al* (2019a) for binary mixtures of sand and biomass, none of the available models (Choi *et al*, 1988; Darton *et al*, 1977; Mori and Wen, 1975) for bubble diameter can predict the bed behavior with a good accuracy. The bed expansion predicted using different models (Agu *et al*, 2019c; Hepbasli, 1998; Werther, 1978) at $U_0/U_{mf} = 2$ show similar trends as the experimental data. Numerically, the Agu *et al* (2019c) model agrees very well with the data from the pellet beds while for the wood chip beds, a good agreement was obtained when the segregated layer at the surface of the bed is excluded in the total bed height measured. The bed expansion model as proposed by Agu *et al* (2019c) is described below:

$$\Delta e = \left[1 - 0.103(U_0 - U_{mf})^{-0.362} \left(\frac{\bar{d}_b}{D} \right)^{-1} \right] - 1 \quad (1)$$

$$\bar{d}_b/D = 0.848 \left(\frac{U_0}{U_{mf}} \right)^{0.66} \left(1 - c \left(\frac{U_0}{U_{mf}} \right)^{a-1} \right)^{0.66};$$

$$Ar > 400 \text{ and } U_0/U_{mf} > c^{1/(1-a)} \quad (2)$$

where $\Delta e = (H_f - H_{mf})/H_{mf}$ is the bed expansion and \bar{d}_b/D is the bubble to bed diameter ratio averaged over the bed height. H_{mf} and H_f are bed heights at the mixture minimum fluidization velocity U_{mf} and at the operating gas velocity U_0 , respectively. Eq. (1) is semi-empirical and was developed based on the mass balance in the solid and gas phases. By using the particle-dependent average bubble diameter shown in Eq. (2) as proposed by Agu *et al* (2018), Eq. (1) predicts the behavior in beds with different particle properties. Ar in Eq. (2) is the Archimedes number, and the equation is applicable for large particles where $Ar > 400$ and for gas velocities U_0 such that $U_0/U_{mf} > c^{1/(1-a)}$ which corresponds to the minimum bubbling condition of the bed, where a and c can be obtained as described in the literature (Agu *et al*, 2018).

For a binary mixture of sand and biomass particles, the mean particle diameter d_m , density ρ_m and sphericity φ_m used for evaluation of \bar{d}_b/D can be obtained from the following, where the subscripts ‘‘m’’, ‘‘b’’ and ‘‘s’’ denote mixture, biomass and sand, respectively, and y_b is the biomass volume fraction in the binary mixture:

$$d_m = \left[\frac{y_b}{\varphi_b d_b} + \frac{(1-y_b)}{\varphi_s d_s} \right]^{-1} \quad (3)$$

$$d_{p, sph, m} = \left[\frac{y_b}{d_b} + \frac{(1-y_b)}{d_s} \right]^{-1} \quad (4)$$

$$\rho_m = y_b \rho_b + (1 - y_b) \rho_s \quad (5)$$

$$\varphi_m = d_m / d_{p, sph, m} \quad (6)$$

3 Results and discussion

Figure 1 compares the predictions of the bed expansion using Eq. (1) with the experimental data obtained at different gas velocities and biomass loads. For the bed of pellets, the mean bed particle properties are based on the total amount of biomass charged in the bed. However, in the wood chip bed, the computation of the mean particle properties is based on 50% of the total biomass charged since approximately 50% of the biomass load lies above the bed surface, preventing the bubble passage (Agu *et al*, 2019a). As can be seen in Figure 1(a), the model results agree very well with the measured data from the wood chip beds, particular with $y_b = 0.2$. For the pellet beds, the model prediction accuracy increases with increasing gas velocity. The inaccuracy in predicting the bed expansion can be associated with the strong interactions between the biomass and sand particles at lower gas velocities. With low gas velocity, a higher amount of sand particles are in the interstices between the biomass particles, thus reducing the bed voidage and expansion. Increasing the gas velocities increases the inter-particle spacing, thereby enhancing the bed expansion.

3.1 Scale up from cold small to hot large bed

To attain a dynamic similarity in the flow of bubbles between a cold small bed and a hot larger bed based on the same biomass properties and volume fraction, the Horio *et al*. (1986) scaling law as described below can be applied.

$$\frac{U_0 - U_{mf}}{\sqrt{gD}}, \frac{U_{mf}}{\sqrt{gD}} \quad (7)$$

The Horio *et al* (1986) law does not impose a restriction on the bed height and on the solid to fluid density ratio. In addition, Glicksman (1988) showed that (7) is equivalent to the scaling law proposed for the viscous limit set for two different beds of the same geometry and negligible solid to fluid density ratio. Implementation of Eq. (7) in this study for scaling up the biomass-sand bed behavior without considering the height and density ratios is based on the fact that bubble growth in a fluidized bed depends on the particle size, fluid properties and bed aspect ratio (Agu *et al*, 2019b). For larger particles or less viscous flow, the bubble growth rate with an increase in gas velocity is high. The bubble size increases with an increase in the bed aspect ratio, particularly for small particle sizes. From (7),

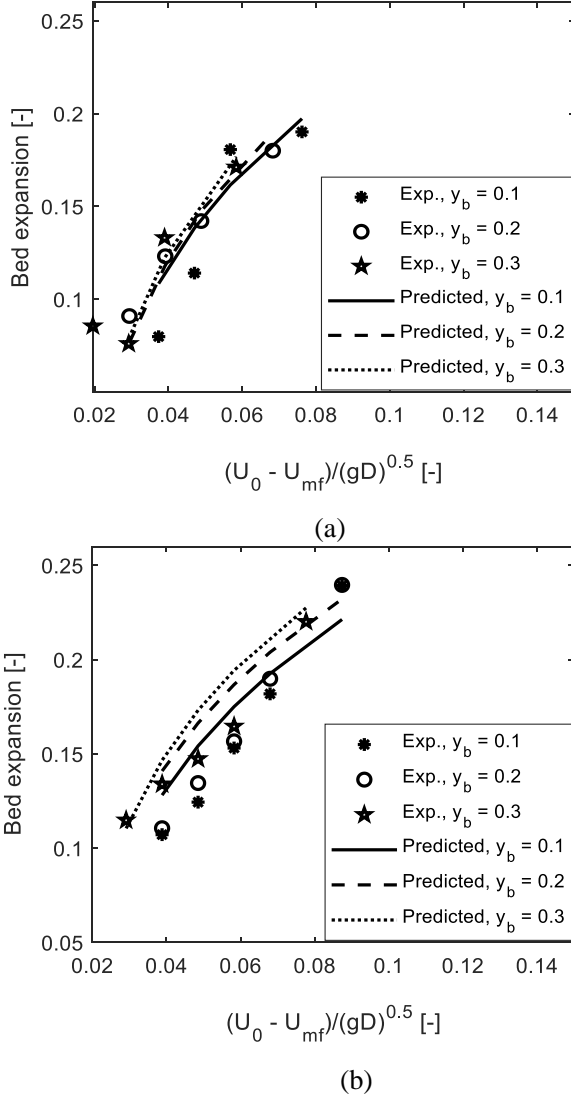


Figure 1. Bed expansion ratio, comparing the predicted results based on Eq. (1) with the experimental data obtained in the cold bed of diameter 10.4 cm at different biomass loads and air velocities as described in Agu *et al.* (2019a); (a) wood chips (b) wood pellets.

$$U_{mf2} = \sqrt{r} U_{mf1} \quad (8)$$

where $r = D_2/D_1$ is the size ratio between the large bed diameter D_2 and the small bed diameter D_1 , and U_{mf2} and U_{mf1} are the corresponding minimum fluidization velocity of the particles in the respective beds. Eq. (8) indicates that particles in the larger bed will have a larger particle size since $U_{mf2} > U_{mf1}$. The particle size for a given value of U_{mf} can be obtained from Eq. (9) as proposed by Wen and Yu (1966a).

$$\frac{\rho_g U_{mf} d_s}{\mu_g} = \left[-33.67 + \sqrt{((33.67)^2 + 0.0408Ar)} \right];$$

$$Ar = \frac{d_s^3 \rho_g (\rho_s - \rho_g) g}{\mu_g^2} \quad (9)$$

If the aspect ratio h_0/D is the same for both beds, the resulting bed height for the larger bed may be too high for industrial application; in addition, the bubble size at

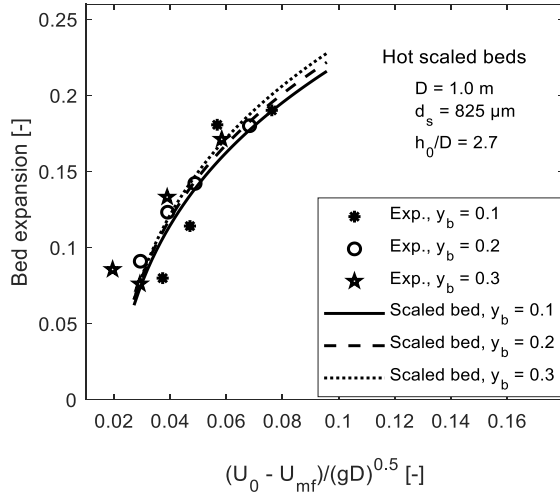
the same $(U_0 - U_{mf})/\sqrt{gD}$ value may be larger. Therefore, for successful application of the Horio *et al.* (1986) scaling law, an output from the large bed that must be correctly matched with that from the smaller bed is required. For the binary mixture of biomass and sand particles, the bed expansion ratio is the output applied in the scaling approach introduced in this study. The bed expansion of the scaled hot larger bed can be simulated from Eq. (1). By comparing the simulated result with the measured data from the cold small bed at the same biomass volume fraction, the appropriate bed aspect ratio can be determined. Since Eq. (1) accurately predicts the behavior in the cold bed of diameter 10.4 cm containing 20% wood chips, this bed is used as the reference here.

Figure 2 compares the bed expansion predicted for the scaled hot large bed with those measured in the cold small bed at different biomass loads. The new bed aspect ratio is evaluated as $h_{0^*}/D_* = 2.7$, where $h_{0^*} = 50$ cm is the initial height of the reference bed (cold bed) and $D_* = 18.5$ cm is the characteristic bed diameter obtained by fitting (1) at the hot bed condition with the data obtained from the cold bed at $y_b = 0.2$. As can be seen, there is a good agreement between the simulated results and the experimental values from the bed with the 20 vol.% wood chips. The similarity also observed for other wood chip beds and those of pellets confirms the validity of the scaling approach.

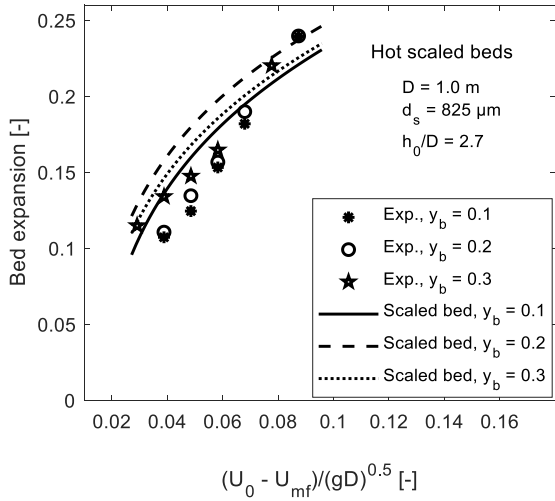
Different dimensionless groups associated with the two beds based on the bed materials are compared in Table 1. In addition to the Froude numbers, $U_{mf}^2/(gD)$, the particle Reynolds numbers, $\rho_g U_{mf} d_s/\mu_g$ and Archimedes numbers, $Ar = d_s^3 \rho_g (\rho_s - \rho_g) g/\mu_g^2$ are closely matched. To achieve the same h_0/d_s ratio, the aspect ratio of the larger bed has to be reduced to $h_0/D = 1.47$, which is possible if only the pure bed material is in the bed. However, by considering that the bed expansion ratio decreases with an amount of biomass in the bed, the height of the large bed needs to be as high as possible to achieve a similar behavior to that in the smaller bed diameter.

3.2 Verifying the scaled bed aspect ratio

To verify the aspect ratio of the scaled large bed, the two different beds (cold at ambient and hot at 800 °C), each containing 20 vol.% wood chips were simulated using Barracuda VR software. Barracuda is a commercially developed platform for implementing the computational particle-fluid dynamics (CPFD) scheme. CPFD is based on the multiphase-particle-in-cell (MP-PIC) concept introduced by Andrew and O'Rourke (1996). In the CPFD scheme, the Euler-Lagrangian modelling approach is applied for fluid volume and particle tracking in gas-solid systems. With the MP-PIC concept, the computational particle represents a large number of particles with similar properties.



(a)



(b)

Figure 2. Bed expansion, comparing the bubbling behavior in the experimental small cold bed (diameter = 10.4 cm, and 293 μm sand particles) with that in the hot scaled bed of diameter 100 cm containing 827 μm sand particles at different biomass loads of (a) wood chips and (b) wood pellets.

The grouping of particles in the CPFDP code makes the simulation faster, thereby increasing its application to industrial systems. Detailed descriptions of the CPFDP model and its numerical scheme can be found in Chen *et al* (2013). The Barracuda software has been successfully used in different studies including investigation of fluid-particle behavior in a binary mixture of bed material and biomass particles (Fotovat *et al*, 2015b) and classification of particles of different sizes and densities (Chladek *et al*, 2018).

In implementing the CPFDP model, the particle properties were the same for both biomass and sand particles as those used in the cold bed experiment (Agu *et al*, 2019a). Both beds ($D = 10.4$ cm and $D = 100$ cm) contain 20 vol% wood chips for each air velocity.

Table 1. Comparison of the dimensionless numbers between the small cold and the large hot beds.

Parameter	Unit	Small bed with 20 vol.% biomass	Large Bed with 20 vol.% biomass
T	$^{\circ}\text{C}$	25	800
p	atm	1	1
D	m	0.104	1.0
ρ_g	kg/m^3	1.18	0.33
μ_g	Pa.s	1.8×10^{-5}	4.3×10^{-5}
d_s	μm	293	827
ρ_s	kg/m^3	2650	2650
U_{mf}	m/s	0.079	0.243
$U_{mf}^2/(gD)$	-	6×10^{-3}	6×10^{-3}
$\rho_g U_{mf} d_s / \mu_g$	-	1.54	1.53
Ar	-	2500	2600
h_0/D	-	4.8	2.7
h_0/d_s	-	1800	3300

The airflow rate at the bottom of the column was assumed uniform across the bed. There are different drag models in the Barracuda library. As the Wen and Yu (1966b) drag model has been used in many applications with good results, this model was also used for all simulations in this study. The default model settings, as given in the Barracuda 17.05 version, were applied.

The geometries are divided into uniform grids. The grid size for the small bed is 13 x 13 x 12.2 mm while that of the larger bed is 9.6 times larger. Relatively coarse grid sizes were used to accommodate the biomass particles (6.87 mm diameter) within each computational cell since in Barracuda it is difficult to capture particles larger than the grid size. Each bed was initialized as a uniform mixture containing 20% wood and 80% sand particles over the bed height. The biomass and sand particles were set as two different particle species to track the individual particle motion and location along the bed. The bottom part of the bed was set as the velocity flow boundary and the top as the pressure boundary. Four different superficial gas velocities were applied in each bed. The simulations were run for 20 s with an initial time step of 0.001 s. The fluid pressure and solid volume fraction at two different positions along the bed axis were used for further analysis.

Figure 3 (a and b) compares the distribution of the solid species (sand and biomass particles) between the two different beds from the initial well mixed state to the final state. For the same biomass volume fraction, the number of biomass particles per unit length is

proportional to the cross-sectional area of the bed. As a result, the scale on the colour bar reflects a higher biomass concentration in the hot bed compared to the cold bed, which is 100 times smaller in cross-sectional area. At the final state, the results show that particle distributions are similar in both beds; most of the biomass particles moved to the bed surface as similarly observed in the experiments. The pressure value at the two indicated positions in the different beds were used as the basis for Figure 3 (c), which shows the variation of pressure drop normalized with the mean value within the last 5 s of the simulations. As can be seen, the fluctuation of pressure drop is similar in both beds, indicating that bubbles rise at approximately the same frequency in the beds. The higher fluctuation amplitude in the hot bed is due to a larger pressure drop over the 70 cm between the measurement positions compared to 15 cm in the case of the cold bed. It should be noted that no scaling rule is applied in choosing the pressure data positions. However, the positions are such that one is within the grid zone (occupying up to 25% of the bed height) and the other within the constant density zone. The variation of mean pressure drop and pressure fluctuations at different gas velocities are shown in Figure 4.

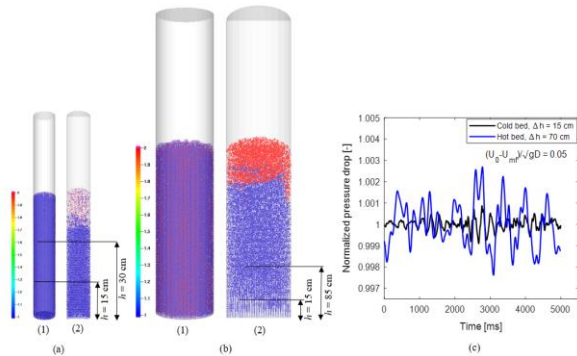


Figure 3. CFPD simulations showing solids distribution (a) cold bed of diameter 10.4 cm, height 50 cm with (1) initial state (2) final state (b) hot (800 °C) scaled bed of diameter 100 cm, height 270 cm with (1) initial state (2) final state (c) pressure drop over the indicated positions within the last-5 s of the simulation time. On the colour scale, “1” indicates sand (blue), “2” indicates biomass (red) and a number in between indicates the particle mixture.

For the same Froude number, $Fr = (U_0 - U_{mf}) / \sqrt{gD}$, Figure 4(a) shows that the pressure drop per unit length is close in the two beds. The approximately constant value of the pressure drop at different gas velocities indicates that the bed is well fluidized. Figure 4(b) gives the relative pressure fluctuation computed as σ/p , where σ is the standard deviation and p is the mean fluid pressure over the last 5 s of the simulation time. The value of σ/p increases with increasing gas velocity due to increasing bubble flux. Up to $Fr = 0.05$, the pressure fluctuations are similar in the two beds, and

also the same at both bed positions. At higher Froude numbers, the difference in the values of σ/p between the two beds (particularly at the upper position) suggests that slugs flow in the beds (although not clear in the simulations). In the cold bed, the experiments show that slugs begin to flow in the upper bed position (observed at 28.7 cm above the distributor) when $Fr = 0.078$. The flow of slugs results in higher fluctuations and bed instability.

With the strong agreement in both the simulated and experimental results, it therefore shows that by using the computed scaled bed aspect ratio, the bubble flow behavior in the hot large bed is similar to that in the cold small bed within the bubble regime.

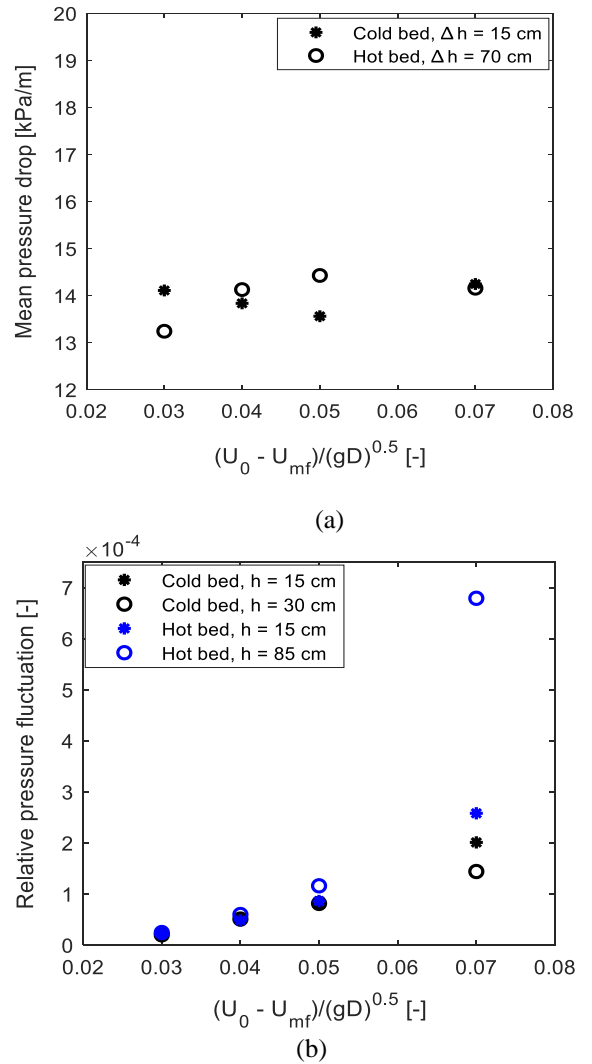


Figure 4. Comparison between cold bed of diameter 10.4 cm, aspect ratio 5 and hot (800 °C) scaled bed of diameter 100 cm, aspect ratio 2.7 at different values of $(U_0 - U_{mf}) / \sqrt{gD}$ (a) mean pressure drop (b) relative fluid pressure fluctuation.

4 Conclusions

In this study, a method for scaling up a cold bed to a hot larger bed was introduced. Using an output such as bed expansion ratio, the proposed scaling approach provides a means of achieving dynamic similarity between two different beds without constraining the fluid and particle properties as well as the bed height, thereby allowing the same properties of fuel particles and bed material type to be used. For illustrations, a 10.4 cm diameter cold bed containing a mixture of biomass and sand particles was scaled up to a large hot bed of 100 cm diameter and the results from CPFD model simulations showed that both beds have clear similarity in bubbling behavior.

Beside bed expansion ratio, other bubble properties including the rise velocity and volumetric bubble flux can also be considered in the scaling approach, provided that such properties are well correlated with the bed geometry and particle properties.

Although the new scaling up approach introduced in this work is verified numerically using commercial simulation software, further verifications using experimental data are necessary.

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